Metastability for an Infinite System of Spiking Neurons



Morgan André

IME-USP

November 27, 2019

▲□▶ ▲□▶ ▲□▶ ▲□▶ ▲□ ● ● ●

Description of the model

The model we are interested in is composed of:

- A countable set I representing the neurons.
- ▶ For each neuron $i \in I$, a set $\mathbb{V}_i \subset I$ of *presynaptic neurons*.
- For each i ∈ I, two point processes (N^{*}_i(t))_{t≥0} and (N[†]_i(t))_{t≥0} representing *spiking times* and *total leak times* respectively.
- For each i ∈ I, a real-valued process (X_i(t))_{t≥0} representing the membrane potential of neuron i.

Spiking times, leaking times and membrane potential

The point process $(N_i^{\dagger}(t))_{t\geq 0}$ is a Poisson process of some rate $\gamma\geq 0$.

The point process $\left(N_i^*(t)\right)_{t\geq 0}$ is characterized by the property that for any $s\leq t$

$$\mathbb{E}\Big(N_i^*(t)-N_i^*(s)|\mathscr{F}_s\Big)=\int_s^t\mathbb{E}\Big(\phi_i(X_i(u))|\mathscr{F}_s\Big)du,$$

▲□▶ ▲□▶ ▲□▶ ▲□▶ ▲□ ● ● ●

Where

$$X_i(t) = \sum_{j \in \mathbb{V}_i} \int_{]L_i(t),t[} dN_j^*(s),$$

and $L_i(t) = \sup\left\{s \le t : N_i^*(\{s\}) + N_i^\dagger(\{s\}) = 1
ight\}$

 $(\phi_i)_{i \in I}$ is a collection of rate functions.

Sub-Critical Metastability

One dimensional lattice with nearest neighbour interaction and hard threshold

(the material presented in this part has been submitted to the journal of statistical physics)

◆□▶ ◆□▶ ◆□▶ ◆□▶ ● ● ●

Spiking times, leaking times and membrane potential

The point process $(N_i^{\dagger}(t))_{t\geq 0}$ is a Poisson process of some rate $\gamma\geq 0$.

The point process $\left(N_i^*(t)\right)_{t\geq 0}$ is characterized by the property that for any $s\leq t$

$$\mathbb{E}\Big(N_i^*(t)-N_i^*(s)|\mathscr{F}_s\Big)=\int_s^t\mathbb{E}\Big(\mathbbm{1}_{X_i(u)>0}|\mathscr{F}_s)\Big)du,$$

Where

$$X_i(t) = \sum_{j \in \mathbb{V}_i} \int_{]L_i(t),t[} dN_j^*(s),$$

and $L_i(t) = \sup \Big\{ s \le t : N_i^*(\{s\}) + N_i^{\dagger}(\{s\}) = 1 \Big\}.$

Spiking times, leaking times and membrane potential

The point process $(N_i^{\dagger}(t))_{t\geq 0}$ is a Poisson process of some rate $\gamma \geq 0$.

The point process $\left(N_i^*(t)\right)_{t\geq 0}$ is characterized by the property that for any $s\leq t$

$$\mathbb{E}\Big(N_i^*(t)-N_i^*(s)|\mathscr{F}_s\Big)=\int_s^t\mathbb{E}\Big(\mathbb{1}_{X_i(u)>0}|\mathscr{F}_s)\Big)du,$$

▲□▶ ▲□▶ ▲□▶ ▲□▶ ▲□ ● ● ●

Where

$$X_{i}(t) = \int_{]L_{i}(t),t[} dN_{i-1}^{*}(s) + \int_{]L_{i}(t),t[} dN_{i+1}^{*}(s),$$

and $L_{i}(t) = \sup \Big\{ s \le t : N_{i}^{*}(\{s\}) + N_{i}^{\dagger}(\{s\}) = 1 \Big\}.$

The choice of $\mathbb{1}_{x>0}$ as the rate function has the important consequence that any given neuron can essentially be in only two states at any given time.

- When X_i(t) = 0 (or equivalently when N^{*}_i(t) has rate 0), we say that neuron i is quiescent.
- When X_i(t) ≥ 1 (or equivalently when N^{*}_i(t) has rate 1), we say that neuron i is active.

Phase transition

Theorem (P. Ferrari et al.)

Suppose that for any $i \in \mathbb{Z}$ we have $X_i(0) \ge 1$. There exists a critical value γ_c for the parameter γ , with $0 < \gamma_c < \infty$, such that for any $i \in \mathbb{Z}$

$$\mathbb{P}\Big(\mathsf{N}^*_i([0,\infty[) < \infty\Big) = 1 \quad \textit{if } \gamma > \gamma_c,$$

and

$$\mathbb{P}\Big(N_i^*([0,\infty[) = \infty\Big) > 0 \quad \text{ if } \gamma < \gamma_c.$$

・ロト ・ 日 ・ ・ 日 ・ ・ 日 ・ ・ つ へ ()

For some $N \in \mathbb{Z}^+$ we set $I_N = \mathbb{Z} \cap [-N, N]$, and we write $(X_i^N(t))_{t\geq 0}$ for the membrane potential of neuron *i* in the version of the model defined on the finite set I_N .

We define the extinction time of this model

$$\sigma_N = \inf \left\{ t \ge 0 : X_i^N(t) = 0 \text{ for all } i \in \mathbb{Z} \cap [-N, N] \right\}.$$

Theorem (Metastability)

If $\gamma < \gamma_{\rm c}$, then we have the following convergence

$$\frac{\sigma_N}{\mathbb{E}(\sigma_N)} \xrightarrow[N \to \infty]{\mathscr{L}} \mathscr{E}(1).$$

うして ふゆう ふほう ふほう うらつ

We consider an important auxiliary process, namely the *spiking* rates process. Denoted $(\xi(t))_{t\geq 0}$ and defined as follows

 $\forall t \geq 0, \ \forall i \in \mathbb{Z}, \quad \xi_i(t) = \mathbb{1}_{X_i(t)>0}.$

This process is an *interacting particle system*. It is a continuous time Markov process taking value in $\{0,1\}^{\mathbb{Z}}$. Each possible state is a doubly infinite sequence of 0 and 1, indicating in which state (quiescent or active) each neuron is.

For some $N \in \mathbb{Z}^+$ we set $I_N = \mathbb{Z} \cap [-N, N]$, and we write $(X_i^N(t))_{t\geq 0}$ for the membrane potential of neuron *i* in the version of the model defined on the finite set I_N .

We define the extinction time of this model

$$\sigma_N = \inf \left\{ t \ge 0 : X_i^N(t) = 0 \text{ for all } i \in \mathbb{Z} \cap [-N, N] \right\}.$$

Theorem (Metastability)

If $\gamma < \gamma_{\rm c}$, then we have the following convergence

$$\frac{\sigma_N}{\mathbb{E}(\sigma_N)} \xrightarrow[N \to \infty]{\mathscr{L}} \mathscr{E}(1).$$

うして ふゆう ふほう ふほう うらつ

For some $N \in \mathbb{Z}^+$ we set $I_N = \mathbb{Z} \cap [-N, N]$, and we write $(\xi_N(t))_{t \ge 0}$ for the spiking rate process restricted to the finite set I_N .

We define the extinction time of this model

$$\sigma_N = \inf \left\{ t \ge 0 : X_i^N(t) = 0 \text{ for all } i \in \mathbb{Z} \cap [-N, N] \right\}.$$

Theorem (Metastability)

If $\gamma < \gamma_{\rm c}$, then we have the following convergence

$$\frac{\sigma_N}{\mathbb{E}(\sigma_N)} \xrightarrow[N \to \infty]{\mathscr{L}} \mathscr{E}(1).$$

For some $N \in \mathbb{Z}^+$ we set $I_N = \mathbb{Z} \cap [-N, N]$, and we write $(\xi_N(t))_{t \ge 0}$ for the spiking rate process restricted to the finite set I_N .

We define the extinction time of the spiking rate process

$$au_{\mathsf{N}} = \inf \Big\{ t \geq 0 : \xi_{\mathsf{N}}(t)_i = 0 ext{ for all } i \in \mathbb{Z} \cap [-\mathsf{N},\mathsf{N}] \Big\}.$$

Theorem (Metastability) If $\gamma < \gamma_c$, then we have the following convergence

$$\frac{\sigma_N}{\mathbb{E}(\sigma_N)} \xrightarrow[N \to \infty]{\mathscr{L}} \mathscr{E}(1).$$

For some $N \in \mathbb{Z}^+$ we set $I_N = \mathbb{Z} \cap [-N, N]$, and we write $(\xi_N(t))_{t \ge 0}$ for the spiking rate process restricted to the finite set I_N .

We define the extinction time of the spiking rate process

$$au_{N} = \inf \Big\{ t \geq 0 : \xi_{N}(t)_{i} = 0 ext{ for all } i \in \mathbb{Z} \cap [-N, N] \Big\}.$$

Theorem (Metastability)

If $\gamma < \gamma_{\rm c}$, then we have the following convergence

$$\frac{\tau_{\mathsf{N}}}{\mathbb{E}(\tau_{\mathsf{N}})} \xrightarrow[N \to \infty]{\mathscr{L}} \mathscr{E}(1).$$

The infinitesimal generator

The generator of the process $(\xi(t))_{t\geq 0}$ is given by

$$\mathscr{L}f(\eta) = \gamma \sum_{i \in \mathbb{Z}} \left(f(\pi_i^{\dagger}(\eta)) - f(\eta) \right) + \sum_{i \in \mathbb{Z}} \eta_i \Big(f(\pi_i(\eta)) - f(\eta) \Big),$$

where the maps are given by

$$\pi_i^{\dagger}(\eta)_j = egin{cases} 0 & ext{if } j=i, \ \eta_j & ext{otherwise}, \end{cases}$$

and

$$\pi_i(\eta)_j = \begin{cases} 0 & \text{if } j = i, \\ \max(\eta_i, \eta_j) & \text{if } j \in \{i - 1, i + 1\}, \\ \eta_j & \text{otherwise.} \end{cases}$$

▲□▶ ▲□▶ ▲□▶ ▲□▶ ▲□ ● ● ●

We can give a graphical construction of the process. For any $i \in \mathbb{N}$ let $(N_i(t))_{t\geq 0}$ and $(N_i^{\dagger}(t))_{t\geq 0}$ be two independent homogeneous Poisson processes with intensity 1 and γ respectively.

Let $(T_{i,n})_{n\geq 0}$ and $(T_{i,n}^{\dagger})_{n\geq 0}$ be their respective jump times.

Consider the time-space diagram $\mathbb{Z} \times \mathbb{R}_+$, and do the following:

- Put a " δ " mark at the point $(i, T_{i,n}^{\dagger})$,
- ▶ Put an arrow pointing from (i, T_{i,n}) to (i + 1, T_{i,n}) and another pointing from (i, T_{i,n}) to (i − 1, T_{i,n}).

うしつ 山 (山) (山) (山) (山) (山) (山) (山)

We say that there is a *valid path* from (i, t) to (j, t') if there is a chain of time segment and arrows leading from (i, t) to (j, t') such that:

- it never cross a " δ " mark,
- when moving upward, we never cross the basis of an arrow.

Then the process $(\xi(t))_{t\geq 0}$ can be defined as follows

$$\xi^{\mathcal{A}}(t) = \{j \in \mathbb{Z}: (i,0) \longrightarrow (j,t) ext{ for some } i \in \mathcal{A}\}$$

・ロト ・ 日 ・ エ ヨ ・ ト ・ 日 ・ う へ つ ・



Figure: In blue all possible valid paths starting from (0,0) up to time t.

◆□▶ ◆□▶ ◆臣▶ ◆臣▶ 三臣 - の々で

The dual process

The process $(\xi(t))_{t\geq 0}$ is additive in the sense that the maps satisfy the following

$$\pi_i(A) = \bigcup_{j \in A} \pi_i(\{j\}) \quad \text{and} \quad \pi_i^{\dagger}(A) = \bigcup_{j \in A} \pi_i^{\dagger}(\{j\})$$

As a consequence it has a dual process, a pure jump Markov process on $\mathscr{P}_f(\mathbb{Z})$ (finite subsets of \mathbb{Z}) denoted $(C(t))_{t\geq 0}$, which generator is given by

$$\mathscr{\tilde{L}}g(F) = \gamma \sum_{i \in F} \Big(g(\tilde{\pi}_i^{\dagger}(F)) - g(F)\Big) + \sum_{i \in F} \eta_i \Big(g(\tilde{\pi}_i(F)) - g(F)\Big),$$

・ロト ・ 日 ・ エ ヨ ・ ト ・ 日 ・ う へ つ ・

The dual maps

For $F \in \mathscr{P}_{f}(\mathbb{Z})$ the dual maps are given by

$$\tilde{\pi}_i^{\dagger}(F) = F \setminus \{i\}$$

 $\quad \text{and} \quad$

$$\tilde{\pi}_i(F) = \bigcup_{j \in F} \tilde{\pi}_i(\{j\})$$

where

$$\tilde{\pi}_i(\{j\}) = \begin{cases} \emptyset & \text{if } j = i, \\ \{i, j\} & \text{if } j \in \{i - 1, i + 1\}, \\ \{j\} & \text{otherwise,} \end{cases}$$

<□▶ <□▶ < □▶ < □▶ < □▶ < □ > ○ < ○

We can give a graphical construction of the dual process as well. Define Poisson processes $(\tilde{N}_i(t))_{t\geq 0}$ and $(\tilde{N}_i^{\dagger}(t))_{t\geq 0}$ with the same parameter as previously, and let $(\tilde{T}_{i,n})_{n\geq 0}$ and $(\tilde{T}_{i,n}^{\dagger})_{n\geq 0}$ be their respective jump times.

Consider the time-space diagram $\mathbb{Z}\times\mathbb{R}_+$ again, and do as for the original process, but reversing the arrows.

A path is said to be *dual-valid* if it satisfies the following constraints:

- it never cross a " δ " mark,
- when moving upward, we never cross the tip of an arrow.

(日) (伊) (日) (日) (日) (0) (0)

The graphical representation of the dual process



Figure: In blue all possible dual-valid paths starting from (0,0) up to time *t*.

▲□▶ ▲圖▶ ▲圖▶ ▲圖▶ _ 圖 _ 釣��

The duality relation

Then we can give an alternative characterization of the dual process as we did for the original process. For any $A \in \mathscr{P}_{f}(\mathbb{Z})$,

$$\mathcal{C}^{\mathcal{A}}(t) = \{j \in \mathbb{Z}: (i,0) \stackrel{ ext{dual}}{\longrightarrow} (j,t) ext{ for some } i \in \mathcal{A} \}.$$

Proposition (Duality)

For any $B \in \mathscr{P}_{f}(\mathbb{Z})$, $A \in \mathscr{P}(\mathbb{Z})$, and $t \geq 0$ we have

$$\mathbb{P}\Big(\xi^{A}(t)\cap B
eq \emptyset\Big)=\mathbb{P}\Big(C^{B}(t)\cap A
eq \emptyset\Big).$$

・ロト ・ 日 ・ エ ヨ ・ ト ・ 日 ・ う へ つ ・

Set monotonicity and stochastic monotonicity

Proposition (Set monotonicity) Let $A \subset B \subset \mathbb{Z}$. Then for any $t \ge 0$ we have

 $\xi^{A}(t) \subset \xi^{B}(t).$

Proposition (Stochastic monotonicity) For any $0 \le s < t$ we have the following

$$\mathbb{P}\left(\xi(s)\in ullet
ight)\geq \mathbb{P}\left(\xi(t)\in ullet
ight).$$

Upper-invariant and lower-invariant measures

Proposition

For any $\gamma > 0$ there exists a probability measures μ_{γ} which is invariant for $(\xi(t))_{t\geq 0}$ and that is such that

$$\mathbb{P}\Big(\xi(t)\in ullet\Big) \underset{t\to\infty}{\longrightarrow} \mu_{\gamma}.$$

The Dirac measure on the "all zero" state, denoted $\delta_{\emptyset},$ is also invariant and we have

$$\mathbb{P}\Big(\xi^{\emptyset}(t)\in ullet\Big) \stackrel{}{\longrightarrow} \delta_{\emptyset}.$$

Moreover if ν is any other invariant measure then $\delta_{\emptyset} \leq \nu \leq \mu_{\gamma}$.

・ロ・・聞・・ヨ・・ 日・ うへの

Upper-invariant and lower-invariant measures for the dual

Proposition

For any $\gamma > 0$ there exists a probability measures $\tilde{\mu}_{\gamma}$ which is invariant for $(C(t))_{t \geq 0}$ and that is such that

$$\mathbb{P}\Big(\mathcal{C}(t)\in oldsymbol{\cdot}\Big) \stackrel{}{\longrightarrow}_{t o\infty} \widetilde{\mu}_{\gamma}.$$

The Dirac measure on the "all zero" state, denoted $\delta_{\emptyset},$ is also invariant and we have

$$\mathbb{P}\Big(C^{\emptyset}(t)\in ullet\Big) \stackrel{}{\longrightarrow} \delta_{\emptyset}.$$

Moreover if ν is any other invariant measure then $\delta_{\emptyset} \leq \nu \leq \tilde{\mu}_{\gamma}$.

・ロト ・西ト ・ヨト ・ヨー うへぐ

Upper-invariant measure in the sub-critical regime

Proposition

When $\gamma < \gamma_c$ we have $\rho_{\gamma} > 0$, and therefore $\mu_{\gamma} \neq \delta_{\emptyset}$.

Proposition When $\gamma < \gamma_c$, we have $\tilde{\mu} \neq \delta_{\emptyset}$.

Proposition

In the sub-critical regime we have $\mu (\eta \equiv 0) = 0$ and $\tilde{\mu} (\eta \equiv 0) = 0$.

・ロト ・ 日 ・ エ ヨ ・ ト ・ 日 ・ う へ つ ・

Space ergodicity of the Upper-invariant measure

Theorem

The measure μ_{γ} is spatially ergodic in the sense that a sequence of random variable $(X_k)_{k \in \mathbb{Z}}$ taking value in $\{0, 1\}$ and such that X_k is distributed like $\mu_{\gamma}(\{\eta : \eta_k = \bullet\})$ would satisfy the following

$$\frac{1}{n+1}\sum_{k=0}^n X_k \xrightarrow[n\to\infty]{a.s.} \rho_{\gamma}.$$

Super-Critical Metastability

One dimensional lattice with nearest neighbour interaction and hard threshold

◆□▶ ◆□▶ ◆□▶ ◆□▶ ● ● ●

(work in progress)

Super-critical metastability

For $\gamma > \gamma_c$ we conjecture that the following convergence holds:

$$\frac{\tau_{\mathsf{N}}}{\mathbb{E}(\tau_{\mathsf{N}})} \xrightarrow[\mathsf{N}\to\infty]{\mathbb{P}} 1.$$

This follows immediately from the two following results (which remain to be proven):

$$\frac{\tau_N}{\log(N)} \xrightarrow[N \to \infty]{\mathbb{P}} C,$$

and

$$\frac{\mathbb{E}(\tau_N)}{\log(N)} \xrightarrow[N \to \infty]{} C.$$

▲□▶ ▲□▶ ▲□▶ ▲□▶ ▲□ ● ● ●

Here C is a positive (and finite) constant.

Simulations

d-dimensional lattice with nearest neighbour interaction and hard threshold

▲□▶ ▲□▶ ▲□▶ ▲□▶ ▲□ ● ● ●

Interaction structure

We did the simulation for $I = \mathbb{Z}$, $I = \mathbb{Z}^2$ and $I = \mathbb{Z}^3$. In each of these cases the structure of the interaction is nearest-neighbors.

In the two-dimensional case the set of presynaptic neurons for neuron i = (i₁, i₂) is

$$\mathbb{V}_i^2 = \{j = (j_1, j_2) \in \mathbb{Z}^2 : |i_1 - j_1| + |i_2 - j_2| = 1\}.$$

► In the three-dimensional case the set of presynaptic neurons for neuron i = (i₁, i₂, i₃) is

$$\mathbb{V}_i^3 = \{j = (j_1, j_2, j_3) \in \mathbb{Z}^3 : |i_1 - j_1| + |i_2 - j_2| + |i_3 - j_3| = 1\}.$$

・ロト ・ 日 ・ エ ヨ ・ ト ・ 日 ・ う へ つ ・

Interaction structure



Figure: Interaction structure in dimension 1 and 2.

▲□▶ ▲圖▶ ▲臣▶ ★臣▶ ―臣 …の�?

Histograms of the extinction time



Figure: Sub-critical simulation for dimension 1, 2 and 3.



Figure: Super-critical simulation for dimension 1, 2 and 3.

◆□> <畳> < Ξ> < Ξ> < □> < □</p>

Simulation for a varying number of neuron



Figure: Fixed (super-critical) γ in the one-dimensional case for a varying number of neurons.

(a)

э

Open problems

<□▶ <□▶ < □▶ < □▶ < □▶ < □ > ○ < ○

Open problems

- In all developments we've evoked the rate function is chosen to be a hard threshold (φ_i(x) = 1_{x>0} for all i ∈ I). It makes the model more manageable mathematically but less realistic biologically. Could we prove the same results for a linear or a sigmoïd rate function?
- In the same lines, the nearest-neighbours structure of the interaction is mathematician-friendly (it allows us to use Harris graphical structure), but could we prove the same results for more complicated graphs (Erdős–Rényi for example)?

Simulation for Erdős-Rényi interaction graph



Figure: Simulation for Erdős–Rényi graph. The parameter of the Erdős–Rényi graph was chosen to be critical, i.e. $p_N = \frac{1}{N}$.

900

ж



Thanks!

◆□▶ ◆□▶ ◆臣▶ ◆臣▶ 臣 の�?